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# Phase transition studied by <sup>7</sup>Li nuclear magnetic resonance in LiXSO<sub>4</sub> (X = K, Rb, Cs and NH<sub>4</sub>) single crystals

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**Abstract.** The temperature dependences of <sup>7</sup>Li nuclear magnetic resonance in LiXSO<sub>4</sub> (X = K, Rb, Cs, and NH<sub>4</sub>) single crystals grown by the slow evaporation method have been investigated by employing a Bruker FT NMR spectrometer. From the experimental data, the nuclear quadrupole constant, the asymmetry parameter and the principal axes of the EFG tensor were determined, and the results were compared with the crystal structure. The temperature dependences of the quadrupole parameters were explained with a single torsional mode of the Li–O bond by the Bayer theory. All the LiO<sub>4</sub> tetrahedra in four different crystals showed torsional motion about the *X*-axis of the EFG tensor. Based on these results, the differences in atomic weight of X in the LiXSO<sub>4</sub> single crystals are responsible for the differences in the torsional angular frequencies.

## 1. Introduction

The increasing number of studies being performed on the physical properties of LiXSO<sub>4</sub> (X = K, Rb, Cs and NH<sub>4</sub>) single crystals are largely owing to its excellent optical quality. These crystals have previously been investigated by means of x-ray diffraction [1–15], Raman spectroscopy [16–28], EPR [29–41], NMR [42–58] and optical properties [59–64]. Also, these materials have ferroelastic properties at low temperatures, as confirmed by stress–strain hysteresis results [65–68] and Aizu classification [69].

In this paper, we reviewed results of the <sup>7</sup>Li nuclear magnetic resonance (NMR) in LiXSO<sub>4</sub> (X = K, Rb, Cs and NH<sub>4</sub>) single crystals grown by the slow evaporation method. The quadrupole coupling constant,  $e^2qQ/h$ , the asymmetry parameter,  $\eta$ , and the direction of the principal axes of the electric field gradient (EFG) tensor of <sup>7</sup>Li (I = 3/2) in LiXSO<sub>4</sub> single crystals were determined as a function of temperature. The temperature dependences of the quadrupole parameters can be discussed with a single torsional frequency of the Li–O ion in the LiO<sub>4</sub> tetrahedron. We found for the first time that the differences in the atomic weight of the alkali ion in each single crystal are affected by the torsional frequency.

#### 2. Crystal structure

These complex sulphates with the composition  $LiXSO_4$  (X = K, Rb, Cs and NH<sub>4</sub>), where X represents an alkali ion, have structures which are based on the  $LiO_4$  tetrahedron. The lithium

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Table 1. The bond lengths for  $LiO_4$  tetrahedra in the  $LiXSO_4$  (X = K, Rb, Cs and NH<sub>4</sub>) single crystals.

LiKSO <sub>4</sub>		LiRbSO <sub>4</sub>		LiCsSO <sub>4</sub>		LiNH <sub>4</sub> SO <sub>4</sub>	
Li–O	1.909 Å	Li–O	1.891 Å	Li–O	1.849 Å	Li–O	1.891 Å
Li–O	1.923 Å	Li–O	1.913 Å	Li–O	1.954 Å	Li–O	1.889 Å
Li–O	1.923 Å	Li–O	1.950 Å	Li–O	1.904 Å	Li–O	1.979 Å
Li–O	1.923 Å	Li–O	1.928 Å	Li–O	1.954 Å	Li–O	1.899 Å

**Table 2.** The structures of  $LiXSO_4$  (X = K, Rb, Cs and NH<sub>4</sub>) single crystals at room temperature.

	LiKSO <sub>4</sub>	LiRbSO <sub>4</sub>	LiCsSO <sub>4</sub>	LiNH <sub>4</sub> SO <sub>4</sub>
Structure Space group Lattice constant	hexagonal P63 a = 5.147 Å	monoclinic P21/n a = 5.288 Å	orthorhombic Pcmn a = 5.456  Å	orthorhombic <i>Pmnn</i> a = 5.28  Å
	c = 8.633  A	b = 9.105  A c = 8.731  Å $\gamma = 90.09^{\circ}$	b = 9.456  A c = 8.820  Å	b = 9.14  A c = 8.786  Å

ion is surrounded by tetrahedrally coordinated oxygens. The bond lengths for  $LiO_4$  tetrahedra are given in table 1. The structures of the  $LiXSO_4$  single crystals are summarized in table 2. The lattice constants and crystal structures are similar to each other, while the phase transition temperatures are absolutely different.

## 3. Experimental procedure

LiXSO<sub>4</sub> (X = K, Rb, Cs and NH<sub>4</sub>) single crystals were grown by slow evaporation from an aqueous solution of equimolar amounts of  $Li_2SO_4 \cdot 4H_2O$  and  $X_2SO_4$  (X = K, Rb, Cs and NH<sub>4</sub>). These single crystals were transparent and colourless. The orientations of the crystal axes were determined by an optical polarizing microscope and the x-ray Laue method. The angular dependence of the NMR spectra was measured for rotations of the crystals around the crystallographic axes *c*, *a* and *b*.

Nuclear magnetic resonance signals of <sup>7</sup>Li in LiXSO<sub>4</sub> single crystals were measured using a Bruker MSL 200 FT NMR spectrometer at the Korea Basic Science Institute in Seoul. The static magnetic field was 4.7 T and the central rf frequency was set at o/2 = 77.777 MHz.

#### 4. Experimental results and analysis

## 4.1. <sup>7</sup>Li NMR in LiKSO<sub>4</sub>

The crystal structure of LiKSO<sub>4</sub> is hexagonal with the space group P63 (C66) and two formula units per cell at room temperature [10]. At lower temperatures, LiKSO<sub>4</sub> undergoes two phase transitions: phase I to phase II at 201 K and phase II to phase III at 190 K. Near 201 K [6], LiKSO<sub>4</sub> transforms from hexagonal phase to either a hexagonal phase [9] or a trigonal phase [17, 59]. Phase III, below 190 K, has been observed to be either orthorhombic [9, 18] or monoclinic [74]. In addition, LiKSO<sub>4</sub> undergoes at least two phase transitions above room temperature: at about 708 and 943 K [70, 71]. The structures of the high-temperature phases are not yet well established. At 708 K the crystal undergoes a structural phase transition, probably to the orthorhombic phase. The second phase transition, most likely to the hexagonal phase, is observed at 943 K. The phase between 708 and 943 K is known to be ferroelastic [72].

## <sup>7</sup>Li NMR in LiXSO<sub>4</sub>



## FREQUENCY (kHz)

Figure 1. The temperature dependence of the <sup>7</sup>Li NMR spectra in an LiKSO<sub>4</sub> crystal.

The <sup>7</sup>Li NMR spectra were measured in the temperature range of 130 to 400 K. The spectrum displayed only three lines for all orientations of the crystal above 190 K, as shown in figure 1. The splitting of <sup>7</sup>Li resonance lines slightly changed in the range above 190 K, which includes the phase transition from I to II. Above phase II, maximum separation of the resonance lines due to the quadrupole interaction was observed when the magnetic field was applied along the *c*-axis of the crystal. This direction was determined to be the *Z*-axis of the EFG tensor. From these NMR results, the quadrupole coupling constant of <sup>7</sup>Li,  $e^2qQ/h = 25$  kHz, and the asymmetry parameter,  $\eta = 0.15$ , were determined at room temperature [50]. These values proved to be non-axially symmetric, consistent with the fact that the coordination of the lithium ion surrounded by oxygen atoms did not form a perfectly regular tetrahedron. Also, because there was no change in the <sup>7</sup>Li NMR spectrum during the phase transition from I to II at 201 K, we could not distinguish between the trigonal and hexagonal space groups proposed for the phase II.

The <sup>7</sup>Li NMR line splits into three lines at the transition point of 190 K as shown in figure 1. The <sup>7</sup>Li NMR spectra below 190 K showed only continuous changes in the quadrupole splittings. From these results, the NMR parameters were obtained. These had the same magnitude of quadrupole coupling constants and asymmetry parameters for resonance lines of each of the three sets, but had different orientations, and were rotated with respect to each other by  $120^{\circ}$  around the *c*-axis. These factors are due to the existence of three kinds of ferroelastic domain. Recently, the structure of the ferroelastic domain is has been discussed in terms of



Figure 2. Temperature dependence of the nuclear quadrupole coupling constant and asymmetry parameter for  $^{7}$ Li NMR in LiKSO<sub>4</sub> crystal.

<sup>7</sup>Li and <sup>39</sup>K NMR results [73]. The principal axes of the <sup>7</sup>Li EFG tensor below 190 K are the same as when it is at room temperature. The principal axes of the Li ion are consistent with the crystallographic axes [51].

The  $e^2 q Q/h$  and  $\eta$  for the <sup>7</sup>Li nucleus in the temperature range of 130–400 K is shown in figure 2. At the transition point of 190 K, the parameters of the <sup>7</sup>Li NMR change abruptly, demonstrating a change in the lithium site symmetry, and the asymmetry parameter  $\eta$  changes from 0.17 above 190 K to 0.44 below 190 K. The quadrupole coupling constant jumps to a value of 34 kHz at the transition point of 190 K. This means that the phase II to III transition is first order. We attempted to explain the temperature dependence of the  $e^2 q Q/h$  and  $\eta$  for <sup>7</sup>Li in terms of a torsional oscillation of the Li–O ion [74, 75]. The values of  $e^2 q Q/h$  and  $\eta$  above and below 190 K were found to decrease almost linearly as a function of increasing temperature;  $e^2 q Q/h$  and  $\eta$  decrease with increasing temperature for this torsional motion about the X-axis. The moment of inertia calculated for the LiKSO<sub>4</sub> crystal structure is  $I_x = 2.73 \times 10^{-45}$  kg m<sup>2</sup>

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## FREQUENCY (kHz)

Figure 3. The temperature dependence of the <sup>7</sup>Li NMR spectra in the LiRbSO<sub>4</sub> crystal.

for the X-axis. The solid lines above and below 190 K in figure 2 are the predicted values of  $e^2 q Q/h$  and  $\eta$  for <sup>7</sup>Li with  $w_x = 3.92 \times 10^{13}$  rad s<sup>-1</sup> and  $w_x = 1.36 \times 10^{13}$  rad s<sup>-1</sup>, respectively.

### 4.2. <sup>7</sup>Li NMR in LiRbSO<sub>4</sub>

Studies on the successive phase transitions in LiRbSO<sub>4</sub> single crystal have been conducted and reported by several groups [76–80]. It has been discovered that LiRbSO<sub>4</sub> crystals undergo successive transitions at 439, 458, 475 and 477 K [76, 77]. These phases have been called I–V in the order of descending temperature. Phases I and V are paraelectric [77]. Phases II and III are ferroelectric and antiferroelectric, respectively [81]. Microscopic observations indicate that phases IV and V are monoclinic, and that the crystal structure is orthorhombic above 458 K [80]. Ferrielectricity has been found in phase IV between 439 and 458 K [77].

The rotation patterns of Li were measured in the three crystallographic planes at room temperature. Because the resonance frequency of the central line is almost constant and the splitting between adjacent lines is equal, the first-order perturbation of  $H_Q$  to  $H_Z$  is sufficient for analysis. By the maximum separation of the resonance line, the *c*-axis of the crystal is found to be the *Z*-axis of the EFG tensor. The quadrupole parameters were determined by a least-squares fit using the experimental data; the quadrupole coupling constant  $e^2 q Q/h = 20.4$  kHz and asymmetry parameter  $\eta = 0$  were determined at room temperature [53].



Figure 4. Temperature dependence of the nuclear quadrupole coupling constant for  $^{7}$ Li NMR in the LiRbSO<sub>4</sub> crystal.

The resonance spectra were measured in the range of 140 to 400 K as shown in figure 3. The line splitting between the central and satellite lines is found to decrease as the temperature increases. Also, the intensity of the resonance line is found to be increased as the temperature increases. The temperature dependence of  $e^2qQ/h$  for <sup>7</sup>Li in the LiRbSO<sub>4</sub> single crystal is shown in figure 4. The value of  $e^2qQ/h$  was found to decrease almost linearly as a function of increasing temperature. The temperature dependence of  $e^2qQ/h$  can be explained by the torsional motion about the X-axis. The moment of inertia,  $I_x = 2.48 \times 10^{-45}$  kg m<sup>2</sup> for the X-axis, was calculated from the bond length of table 1. The inertia moment is calculated by the nearest four oxygen ions. The solid lines in figure 4 are the predicted values of  $e^2qQ/h$  for <sup>7</sup>Li with  $w_x = 8.30 \times 10^{12}$  rad s<sup>-1</sup>. Therefore, the Bayer–Wang theory can satisfactorily explain our data for the temperature range 140–400 K.

## 4.3. <sup>7</sup>Li NMR in LiCsSO<sub>4</sub>

LiCsSO<sub>4</sub> single crystals undergo a transition from the paraelastic phase with orthorhombic structure at room temperature to the ferroelastic phase with monoclinic structure below  $T_c = 202$  K [9, 82].

The <sup>7</sup>Li NMR spectra were measured in the temperature range of 140 to 400 K, as shown in figure 5. Above  $T_c$ , the <sup>7</sup>Li NMR spectrum showed three lines for all orientations of the magnetic field, while two sets of <sup>7</sup>Li NMR spectra were recorded below  $T_c$ . From the rotation patterns of <sup>7</sup>Li NMR spectra in the three crystallographic planes, the nuclear quadrupole coupling constant and asymmetry parameter of <sup>7</sup>Li in an LiCsSO<sub>4</sub> crystal were determined as 22.35 kHz and 0.68 at room temperature. The satellite resonance lines did not show the extrema along the crystallographic axis. Based on these experimental results, we can conclude that the principal axes of the Li ion are not parallel to the crystallographic axes. The EFG tensor of <sup>7</sup>Li was found to be asymmetric. The direction of the principal EFG tensor at 300 K is represented with the Eulerian angles  $\Phi = 90^\circ$ ,  $\Theta = 20^\circ$  and  $\Psi = 0^\circ$  [54].

At the transition point of 200 K, the <sup>7</sup>Li NMR line splits into two sets. The rotation patterns of the <sup>7</sup>Li NMR spectra are measured in the ab-, bc- and ca-plane below 200 K. The obtained results can be explained by the existence of two kinds of ferroelastic domain, rotated with respect to each other by 120° around the c-axis. Below 200 K, the <sup>7</sup>Li NMR spectra showed



## FREQUENCY (kHz)

Figure 5. The temperature dependence of the <sup>7</sup>Li NMR spectra in the LiCsSO<sub>4</sub> crystal.

only continuous quantitative changes in the quadrupole splittings, without any other abrupt changes. The <sup>7</sup>Li NMR spectra demonstrated the occurrence of a phase transition at 200 K, which is connected with a lowering of the Li<sup>+</sup> site symmetry and formation of two kinds of ferroelastic domain. The experimental data show that the principal axes of the <sup>7</sup>Li EFG tensor below  $T_c$  are not the same as those above  $T_c$ . This means that the phase transition is second order. The principal X-, Y- and Z-axes are found to lie along the Eulerian angles  $\Phi = 70^\circ$ ,  $\Theta = 15^\circ$  and  $\Psi = 35^\circ$ . The nuclear quadrupole coupling constant and asymmetry parameter of <sup>7</sup>Li in an LiCsSO<sub>4</sub> crystal were determined to be 26.90 kHz and 0.79 at 180 K [54].

The temperature dependences of  $e^2 q Q/h$  and  $\eta$  in the temperature range of 140–400 K are shown in figure 6. For the torsional motion about the *X*-axis, both  $e^2 q Q/h$  and  $\eta$  decrease with increasing temperature. The Mathematica package was used to simulate the variations of  $e^2 q Q/h$  and  $\eta$ . We controlled sensitivity variations of  $w_x$  for the slopes of the  $e^2 q Q/h$  and  $\eta$  obtained from our NMR experimental results. The straight line in figure 6 is the prediction of  $e^2 q Q/h$  and  $\eta$  for <sup>7</sup>Li with  $I_x = 1.67 \times 10^{-45}$  kg m<sup>2</sup> and  $w_x = 4.50 \times 10^{12}$  rad s<sup>-1</sup>.

## 4.4. <sup>7</sup>Li NMR in LiNH<sub>4</sub>SO<sub>4</sub>

LiNH<sub>4</sub>SO<sub>4</sub> single crystals undergo two phase transitions [42, 83, 84] with  $T_{c1}$  at 459 K and  $T_{c2}$  at approximately 283 K. These take the crystal from the orthorhombic high temperature phase of space group  $D_{2h}^{16}$  (z = 4) [85] to another orthorhombic phase, space group  $C_{2v}^{9}$  [86],



Figure 6. Temperature dependence of the nuclear quadrupole coupling constant and asymmetry parameter for  $^{7}$ Li NMR in the LiCsSO<sub>4</sub> crystal.

and then transform it further to a monoclinic phase, space group  $C_{2h}^5$  [87]. Three phases of LiNH<sub>4</sub>SO<sub>4</sub> can be distinguished: phase I for T > 459 K, phase II for 283 K < T < 459 K, and phase III for T < 283 K. It has been established that the phase transition from phase II to III is of the first order, while that from phase I to II is of the second order [43].

In the temperature range of 170 to 400 K, the <sup>7</sup>Li NMR spectra were measured as shown in figure 7. A <sup>7</sup>Li NMR spectrum consists of only three lines for all orientations of the crystal above  $T_{c2}$  (=283 K). From these NMR results, the quadrupole coupling constant,  $e^2q Q/h = 25$  kHz, and asymmetry parameter,  $\eta = 0.22$ , were obtained at room temperature. The EFG tensor of <sup>7</sup>Li was found to be non-axially symmetric. At room temperature, the satellite resonance lines showed neither maximum nor minimum separation along the crystallographic axes. Thus, we can conclude that the principal axes of the EFG tensors at room temperature are represented with the Eulerian angles  $\Phi = 80^\circ$ ,  $\Theta = 80^\circ$  and  $\Psi = 8.5^\circ$  [58].

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Figure 7. The temperature dependence of the <sup>7</sup>Li NMR spectra in the LiNH<sub>4</sub>SO<sub>4</sub> crystal.

The three-line structure is a result of the quadrupole interaction of the <sup>7</sup>Li (I = 3/2) nucleus. However, two different groups of Li resonances were recorded in phase III, originating from Li(I) and Li(II), respectively. These two signals are associated with two physically inequivalent positions of lithium atoms in the unit cell. The rotation pattern of Li(I) and Li(II) was measured in three crystallographic planes below 287 K. Here, the signal for Li(II) showed very strong intensity, while the Li(I) showed very weak intensity. Two different Li resonance groups with different magnitudes of quadrupole splitting were analysed. At 200 K, the quadrupole coupling constant and asymmetry parameter were determined as 25.5 kHz and 0.33 for Li(I), 34 kHz and 0.88 for Li(II), respectively. The principal axes of the EFG tensor at low temperatures were not the same as those obtained above  $T_c$ . The principal axes X, Y and Z of the EFG tensor lie along the Eulerian angles  $\Phi = 90^{\circ}$ ,  $\Theta = 90^{\circ}$ ,  $\Psi = 0^{\circ}$  and  $\Phi = 70^{\circ}$ ,  $\Theta = 87^{\circ}$ ,  $\Psi = 79^{\circ}$ , for Li(I) and Li(II), respectively. The EFG tensors of Li(I) and Li(II) are both asymmetric, and the orientations of the principal axes of the EFG tensors also do not coincide for Li(I) and Li(II) [57].

The temperature dependences of  $e^2 q Q/h$  and  $\eta$  are shown in figure 8. In the temperature range 170 to 400 K, which includes the II–III phase transition, there is an abrupt change in the parameters of <sup>7</sup>Li NMR. This means that the phase transition is first order. At the transition point of 287 K, the <sup>7</sup>Li NMR line splits into two sets. Below 287 K, the quadrupole parameter of Li(I) slowly increases as the temperature decreases, while that of Li(II) was found to drastically



Figure 8. Temperature dependence of the nuclear quadrupole coupling constant and asymmetry parameter for  $^{7}$ Li NMR in the LiNH<sub>4</sub>SO<sub>4</sub> crystal.

increase with the temperature. The moment of inertia calculated from the crystal structure is  $I_x = 2.56 \times 10^{-45}$  kg m<sup>2</sup> for the X-axis. The solid lines above 287 K are the predicted values of  $e^2 q Q/h$  and  $\eta$  with  $w_x = 1.60 \times 10^{13}$  rad s<sup>-1</sup> for Li. Below 287 K, the values of  $e^2 q Q/h$  and  $\eta$  are  $w_x = 1.60 \times 10^{13}$  rad s<sup>-1</sup> and  $w_x = 1.03 \times 10^{13}$  rad s<sup>-1</sup> for Li(I) and Li(II), respectively.

## 5. Discussion and conclusion

The <sup>7</sup>Li NMR in the single crystal form of LiXSO<sub>4</sub> (X = K, Rb, Cs and NH<sub>4</sub>) grown by the slow evaporation method has been investigated by employing a Bruker FT NMR spectrometer. From the experimental data, the quadrupole coupling constant, asymmetry parameter and the direction of EFG tensor were determined as functions of temperature. The nuclear electric quadrupole interaction of the <sup>7</sup>Li nucleus having the nuclear spin I = 3/2 provides information

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about the electric field gradient produced by ions surrounding the resonant nucleus. Thus, the quadrupole coupling constant and asymmetry parameter of <sup>7</sup>Li reveal the configuration of ionic charges around the Li<sup>+</sup>. The  $e^2qQ/h$  values for the <sup>7</sup>Li nucleus in the LiXSO<sub>4</sub> crystal system were found to be similar to each other. This similarity is consistent with the fact that LiXSO<sub>4</sub> single crystals have isomorphous structures. The differences in the field gradient in the case of each crystal were observed at different positions of the Li ions. The splitting of <sup>7</sup>Li NMR spectra at  $T_c$  demonstrates the occurrence of a phase transition which is connected with a lowering of the Li<sup>+</sup> site symmetry and also indicates the formation of ferroelastic domains.

From the NMR of <sup>7</sup>Li in LiXSO<sub>4</sub>, we can observe the mechanism of phase transitions. In the case of LiKSO<sub>4</sub> and LiNH<sub>4</sub>SO<sub>4</sub> crystals, the NMR parameter shows a discontinuity at  $T_c$ . This means that the phase transition is a first-order transition. Also, the parameters of the LiCsSO<sub>4</sub> crystal show continuity at  $T_c$ , and undergo a second-order phase transition.

Meng and Cao [46] have made a <sup>7</sup>Li NMR analysis of LiKSO<sub>4</sub> at room temperature. According to their results, the rotation pattern in the *ab*-plane was angle independent, while the rotation patterns in the *ac*- and *bc*-plane were angle dependent. The above result indicates that the electric field gradient tensors of <sup>7</sup>Li were axially symmetric. However, the <sup>7</sup>Li NMR parameter in LiKSO<sub>4</sub> crystals studied by our results were non-axially symmetric. This result is not consistent with the result of other groups. Consequently, LiKSO<sub>4</sub> single crystals may have different crystal structures according to the conditions of crystal growth.

The proposed torsional angular frequencies in  $LiXSO_4$  single crystals were determined from the temperature dependent quadrupole parameters. Based on these results, the differences in atomic weight of X in the  $LiXSO_4$  single crystals are responsible for the differences in the torsional angular frequencies.

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